

SM315 Lecture Notes  
 Vibrating Circular Membrane  
 Homework: (TBD)

## 1. Wave Equation in 2D Polar Coordinates

a. PDE:  $\frac{\partial^2 u}{\partial t^2} = \omega^2 \nabla^2 u = \omega^2 \left( \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} \right)$  where  $\begin{cases} 0 \leq r \leq a \\ -\pi \leq \theta \leq \pi \\ t \geq 0 \end{cases}$

b. BC:  $\begin{cases} u(a, \theta, t) = 0 \\ |u(0, \theta, t)| < \infty \end{cases}$  and  $\begin{cases} u(r, -\pi, t) = u(r, \pi, t) \\ \frac{\partial u}{\partial \theta}(r, -\pi, t) = \frac{\partial u}{\partial \theta}(r, \pi, t) \end{cases}$

- $u(a, \theta, t) = 0$  implies that outside edge of circle is pinned down (i.e. like a drum head).
- $|u(0, \theta, t)| < \infty$  implies that the displacement must be finite.

c. IC:  $\begin{cases} u(r, \theta, 0) = f(r, \theta) \\ \frac{\partial u}{\partial t}(r, \theta, 0) = g(r, \theta) \end{cases}$

- BC's in theta are due to the periodic nature of the solution

## 2. Separate the Variables

- a. Expand the PDE by applying product rule to  $\frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u}{\partial r} \right)$  term, i.e.:

$$\frac{\partial^2 u}{\partial t^2} = \omega^2 \left( \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} \right)$$

- b.  $u(r, \theta, t) = R(r)\Theta(\theta)T(t)$  or in shorthand  $u = R\Theta T$

- c. PDE becomes:  $R\Theta T'' = \omega^2 R''\Theta T + \frac{\omega^2}{r} R'\Theta T + \omega^2 R\Theta'' T$

- d. Now divide by  $\omega^2 R\Theta T \rightarrow$  PDE becomes:

$$\frac{1}{\omega^2} \frac{T''}{T} = \frac{R''}{R} + \frac{1}{r} \frac{R'}{R} + \frac{\Theta''}{\Theta}$$

- e. Note that the right hand side is a function of  $t$  only, and the left hand side is a function of  $r$  and  $\theta$  only. The only way for this equality to exist is for both sides to be equal to the same constant  $-\lambda$ , i.e.:

$$\frac{1}{\omega^2} \frac{T''}{T} = \frac{R''}{R} + \frac{1}{r} \frac{R'}{R} + \frac{\Theta''}{r^2 \Theta} = -\lambda =$$

- f. This generates the familiar ODE for  $T$ :  $T'' + \lambda \omega^2 T = 0$   
 g. In order to evaluate the eigenvalue, boundary conditions must be applied. But  $r$  and  $\theta$  are not yet separated.

- h. In order to separate  $r$  from  $\theta$  we must multiply  $\frac{R''}{R} + \frac{1}{r} \frac{R'}{R} + \frac{\Theta''}{r^2 \Theta} = -\lambda$  by  $r^2$ , i.e.

$$\frac{r^2 R''}{R} + \frac{rR'}{R} + \frac{\Theta''}{\Theta} = -\lambda$$

- i. So we proceed in this manner separating  $r$  from  $\theta$ :

$$\frac{r^2 R''}{R} + \frac{rR'}{R} + \lambda r^2 = -\frac{\Theta''}{\Theta} = \mu$$

- The reason for using  $\mu$  vice  $-\mu$  is evident in the next step

- j. This generates the ODE for  $\theta$ :  $\Theta'' + \mu\Theta = 0$

- k. And finally the ODE for  $R$ :  $r^2 R'' + rR' + (\lambda r^2 - \mu)R = 0$

- l. In Summary, the three ODEs are:

$$\begin{cases} T'' + \lambda \omega^2 T = 0 \\ \Theta'' + \mu \Theta = 0 \\ r^2 R'' + rR' + (\lambda r^2 - \mu)R = 0 \end{cases}$$

### 3. Use the BC's in $\theta$ to Evaluate the Eigenvalue $\mu$

a. We have solved this boundary value problem before when we solved the problem in a circular ring (See Table 2.4.1 on Page 69):

- $\mu_m = \left(\frac{m\pi}{L}\right)^2$  where  $m = 0, 1, 2, \dots$ . We consider that  $L = \pi$  since  $-\pi \leq \theta \leq \pi$  and therefore  $\mu_m = m^2$
- Thus  $\Theta'' + m^2\Theta = 0 \rightarrow \Theta = A_m \sin(m\theta) + B_m \cos(m\theta)$
- At this point we can do nothing more with the  $\theta$  BC

### 4. Use the BC's in $r$ to Evaluate the Eigenvalue $\lambda$

a. The differential equation in  $r$ , although unfamiliar to us, is actually a very well known equation (after some manipulation)

b. We let

$$\begin{cases} z = \sqrt{\lambda}r \rightarrow r = z/\sqrt{\lambda} \rightarrow \\ \frac{dR}{dr} = \frac{dR}{dz} \frac{dz}{dr} = \sqrt{\lambda} \frac{dR}{dz} \rightarrow \\ \frac{d^2R}{dr^2} = \frac{d}{dr} \frac{dR}{dr} = \frac{d}{dr} \left[ \sqrt{\lambda} \frac{dR}{dz} \right] = \lambda \frac{d^2R}{dz^2} \end{cases}$$

c. therefore:

$$\begin{cases} r^2 \frac{d^2}{dr^2} R(r) + r \frac{d}{dr} R(r) + (\lambda r^2 - m^2) R(r) = 0 \rightarrow \\ \frac{z^2}{\lambda} \frac{d^2}{dz^2} R(z(r)) + \frac{z}{\sqrt{\lambda}} \frac{d}{dz} R(z(r)) + (z^2 - m^2) R(z(r)) = 0 \end{cases}$$

d. Now apply chain rule:

$$\begin{cases} \frac{d}{dz} R(z(r)) = \frac{dR}{dz} \frac{dz}{dr} = \frac{dR}{dz} \frac{d}{dz} (\sqrt{\lambda}r) = \sqrt{\lambda} \frac{dR}{dz} \\ \frac{d^2}{dr^2} = \frac{d}{dr} \left( \sqrt{\lambda} \frac{dR}{dz} \right) = \sqrt{\lambda} \frac{d^2R}{dz^2} \frac{dz}{dr} = \sqrt{\lambda} \frac{d^2R}{dz^2} \frac{d}{dr} (\sqrt{\lambda}r) = \lambda \frac{d^2R}{dz^2} \end{cases}$$

e. Substituting this in b yields:

$$z^2 \frac{d^2R}{dz^2} + z \frac{dR}{dz} + (z^2 - m^2)R = 0$$

f. This is known as **Bessel's Equation of Order  $m$**

*Bessel's Differential Equation* which arises in numerous problems, especially in polar and cylindrical coordinates. **Bessel's Differential Equation** is defined as:

$$x^2 y'' + xy' + (x^2 - \#^2)y = 0$$

where  $\#$  is a non-negative real number. The solutions of this equation are called **Bessel Functions** of order  $\#$ . Although the order  $\#$  can be any real number, the scope of this section is limited to *non-negative integers*, i.e.,  $\# = 0, 1, 2, 3, \dots$ , unless specified otherwise.

## 5. Notes on Bessel

- a. Bessel worked out a method of mathematical analysis involving what is now known as the Bessel function. He introduced this in 1817 in his study of a problem of [Kepler](#) of determining the motion of three bodies moving under mutual gravitation.
- b. Bessel functions appear as coefficients in the series expansion of the indirect perturbation of a planet, that is the motion caused by the motion of the Sun caused by the perturbing body. In 1824 he developed Bessel functions more fully in a study of planetary perturbations and published a treatise on them in Berlin. It was not the first time that special cases of the functions had appeared, [Jacob Bernoulli](#), [Daniel Bernoulli](#), [Euler](#) and [Lagrange](#) having studied special cases of them earlier. In fact it was probably [Lagrange](#)'s work on elliptical orbits that first suggested to Bessel to work on the Bessel functions.

